

Explicit fundamental solutions of some second order differential operators on Heisenberg groups

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Abstract

Let p, q, n be natural numbers such that $p + q = n$. Let \mathbb{F} be either \mathbb{C} , the complex numbers field, or \mathbb{H} , the quaternionic division algebra. We consider the Heisenberg group $N(p, q, \mathbb{F})$ defined as $N(p, q, \mathbb{F}) = \mathbb{F}^n \times \mathfrak{Im}\mathbb{F}$, with group law given by

$$(v, \zeta)(v', \zeta') = \left(v + v', \zeta + \zeta' - \frac{1}{2}\mathfrak{Im}B(v, v') \right),$$

where $B(v, w) = \sum_{j=1}^p v_j \overline{w_j} - \sum_{j=p+1}^n v_j \overline{w_j}$. Let $U(p, q, \mathbb{F})$ be the group of $n \times n$ matrices with coefficients in \mathbb{F} that leave invariant the form B . In this work we compute explicit fundamental solutions of some second order differential operators on $N(p, q, \mathbb{F})$ which are canonically associated to the action of $U(p, q, \mathbb{F})$.

1 Introduction and Preliminaries

Let p, q, n be natural numbers such that $p + q = n$. Let \mathbb{F} be either \mathbb{C} , the complex numbers field, or \mathbb{H} , the quaternionic division algebra. We consider the Heisenberg group $N(p, q, \mathbb{F})$ defined as $N(p, q, \mathbb{F}) = \mathbb{F}^n \times \mathfrak{Im}\mathbb{F}$, with group law given by

$$(v, \zeta)(v', \zeta') = \left(v + v', \zeta + \zeta' - \frac{1}{2}\mathfrak{Im}B(v, v') \right),$$

where $B(v, w) = \sum_{j=1}^p v_j \overline{w_j} - \sum_{j=p+1}^n v_j \overline{w_j}$. The associated Lie algebra is $\eta(p, q, \mathbb{F}) = \mathbb{F}^n \oplus \mathfrak{Im}\mathbb{F}$, with Lie bracket given by

$$[(v, \zeta), (v', \zeta')] = (0, -\mathfrak{Im}B(v, v')).$$

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Let $\mathcal{U}(\eta(p, q, \mathbb{F}))$ be the universal enveloping algebra of $\eta(p, q, \mathbb{F})$ which we identify with the algebra of left invariant differential operators. Let $U(p, q, \mathbb{H})$ be the group of $n \times n$ matrices with coefficients in \mathbb{F} that leave invariant the form B . Then $U(p, q, \mathbb{F})$ acts by automorphism on $\eta(p, q, \mathbb{F})$ by

$$g \cdot (v, \zeta) = (gv, \zeta).$$

We denote by $\mathcal{U}(\eta(p, q, \mathbb{F}))^{U(p, q, \mathbb{F})}$ the subalgebra of $\mathcal{U}(\eta(p, q, \mathbb{F}))$ of the left invariant differential operators which commute with this action. It is known that this subalgebra is generated by two operators: L and U , and a family of tempered joint eigendistributions is computed explicitly (see for example [D-M], [G-S(1)], [V]).

More precisely, if $\mathbb{F} = \mathbb{C}$ and $\{X_1, \dots, X_n, Y_1, \dots, Y_n, U\}$ denotes the standard basis of the Heisenberg Lie algebra with $[X_i, Y_j] = \delta_{ij}U$ and all the other brackets are zero, then

$$L = \sum_{j=1}^p X_j^2 + Y_j^2 - \sum_{j=p+1}^n X_j^2 + Y_j^2.$$

For $\lambda \in \mathbb{R}$, $\lambda \neq 0$ and $k \in \mathbb{Z}$, $S_{\lambda, k}$ is a $U(p, q)$ -invariant tempered distribution satisfying

$$\begin{aligned} LS_{\lambda, k} &= -|\lambda|(2k + p - q)S_{\lambda, k}, \\ iUS_{\lambda, k} &= \lambda S_{\lambda, k}. \end{aligned}$$

This family provides us an *inversion formula*: for all f in the Schwartz space on the Heisenberg group, we have that

$$f(z, t) = \sum_{k \in \mathbb{Z}} \int_{-\infty}^{\infty} f * S_{\lambda, k} |\lambda|^n d\lambda, \quad (z, t) \in N(p, q, \mathbb{C}). \quad (1.1)$$

If $\mathbb{F} = \mathbb{H}$ we take $\{X_1^0, X_1^1, X_1^2, X_1^3, \dots, X_n^0, X_n^1, X_n^2, X_n^3, Z_1, Z_2, Z_3\}$ the canonical basis for the Lie algebra, where Z_1, Z_2, Z_3 generate the center of $\eta(p, q, \mathbb{H})$. Here,

$$\begin{aligned} L &= \sum_{r=1}^p \sum_{l=0}^3 (X_r^l)^2 - \sum_{r=p+1}^n \sum_{l=0}^3 (X_r^l)^2, \quad \text{and} \\ U &= \sum_{l=1}^3 Z_l^2. \end{aligned}$$

There also exists a family of $U(p, q, \mathbb{H})$ -invariant tempered distributions $\varphi_{w, k}$, $w \in \mathbb{R}^3$ y $k \in \mathbb{Z}$, such that each one of them is a joint eigendistribution of L and U :

$$\begin{aligned} L\varphi_{w, k} &= -|w|(2k + 2(p - q))\varphi_{w, k}, \\ U\varphi_{w, k} &= -\lambda^2\varphi_{w, k}; \end{aligned}$$

in this case this family also provides an inversion formula: for all $f \in \mathcal{S}(N(p, q, \mathbb{H}))$ we have that

$$f(\alpha, z) = \sum_{k \in \mathbb{Z}} \int_{\mathbb{R}^3} (f * \varphi_{w, k})(\alpha, z) |w|^{2n} dw, \quad (\alpha, z) \in N(p, q, \mathbb{H}). \quad (1.2)$$

The aim of this work is to explicitly compute a fundamental solution in the classical case for the operator $\mathcal{L}_\alpha = L + i\alpha U$, where α is a complex number; and in the quaternionic case for the operator L . Recall that a fundamental solution for the differential operator \mathcal{L} is a tempered distribution Φ such

that for all f in the Schwartz class, we have that $\mathcal{L}(f * \Phi) = (\mathcal{L}f) * \Phi = f * \mathcal{L}(\Phi) = f$. So if we define the operator K as $Kf = f * \Phi$, then $K \circ \mathcal{L}f = \mathcal{L} \circ Kf = f$.

From the inversion formula (1.1) it is natural to propose as a fundamental solution of \mathcal{L}_α

$$\langle \Phi_\alpha, f \rangle = \sum_{k \in \mathbb{Z}} \int_{-\infty}^{\infty} \frac{1}{-|\lambda|(2k + p - q - \alpha \operatorname{sgn} \lambda)} \langle S_{\lambda,k}, f \rangle |\lambda|^n d\lambda, \quad (1.3)$$

for $f \in \mathcal{S}(N(p, q, \mathbb{C}))$; and from (1.2) we propose as a fundamental solution of L

$$\langle \Phi, f \rangle = \sum_{k \in \mathbb{Z}_{\mathbb{R}^3}} \int \frac{1}{-|w|(2k + 2(p - q))} \langle \varphi_{w,k}, f \rangle |w|^{2n} dw, \quad (1.4)$$

for $f \in \mathcal{S}(N(p, q, \mathbb{H}))$.

We remark that for $q = 0$, $\mathbb{F} = \mathbb{C}$ we recover the fundamental solution for the operator \mathcal{L}_α given in [F-S], and for $q = 0$, $\mathbb{F} = \mathbb{H}$ we recover Kaplan's fundamental solution for the operator L given in [K]. The case $q \neq 0$, $\alpha = 0$ was obtained in [G-S(2)].

The expression of Φ_α is obtained in theorem 2.9, and for the computation we follow the method used in [G-S(2)]. In the quaternionic case, Φ is given in theorem 3.1, and for its computation we use the Radon transform in order to reduce this case to the classical one.

To describe both families of eigendistributions $\{S_{\lambda,k}\}$ and $\{\varphi_{w,k}\}$ we need to adapt a result by Tengstrand in [T]. We describe the elements for the case $\mathbb{F} = \mathbb{C}$, the other one is similar. First of all we take bipolar coordinates on \mathbb{C}^n for $(x_1, y_1, \dots, x_n, y_n)$ we set $\tau = \sum_{j=1}^p (x_j^2 + y_j^2) - \sum_{j=p+1}^n (x_j^2 + y_j^2)$, $\rho = \sum_{j=1}^n (x_j^2 + y_j^2)$, $u = (x_1, y_1, \dots, x_p, y_p)$, $v = (x_{p+1}, y_{p+1}, \dots, x_n, y_n)$. Hence $u = (\frac{\rho+\tau}{2})^{\frac{1}{2}} \omega_u$, with $\omega_u \in S^{2p-1}$ and $v = (\frac{\rho-\tau}{2})^{\frac{1}{2}} \omega_v$, with $\omega_v \in S^{2q-1}$. It is easy to see by changing variables that

$$\begin{aligned} \int_{\mathbb{C}^n} f(z) dz &= \int_{-\infty}^{\infty} \int_{\rho > |\tau|} \int_{S^{2p-q} \times S^{2q-1}} f \left(\left(\frac{\rho+\tau}{2} \right)^{\frac{1}{2}} \omega_u, \left(\frac{\rho-\tau}{2} \right)^{\frac{1}{2}} \omega_v \right) d\omega_u d\omega_v \times \\ &\quad \times (\rho + \tau)^{p-1} (\rho - \tau)^{q-1} d\rho d\tau. \end{aligned}$$

Then we define for $f \in \mathcal{S}(\mathbb{R}^{2n})$

$$Mf(\rho, \tau) = \int_{S^{2p-1} \times S^{2q-1}} f \left(\left(\frac{\rho+\tau}{2} \right)^{\frac{1}{2}} \omega_u, \left(\frac{\rho-\tau}{2} \right)^{\frac{1}{2}} \omega_v \right) d\omega_u d\omega_v,$$

and also

$$Nf(\tau) = \int_{|\tau|}^{\infty} Mf(\rho, \tau) (\rho + \tau)^{p-1} (\rho - \tau)^{q-1} d\rho.$$

Let us now define the space \mathcal{H}_n of the functions $\varphi : \mathbb{R} \rightarrow \mathbb{C}$ such that $\varphi(\tau) = \varphi_1(\tau) + \tau^{n-1} \varphi_2(\tau) H(\tau)$, for $\varphi_1, \varphi_2 \in \mathcal{S}(\mathbb{R})$, where H denotes the Heaviside function. In [T] it is proved that \mathcal{H}_n with a suitable topology is a Fréchet space, and following the same lines we can see that the linear maps $N : \mathcal{S}(\mathbb{R}^{2n} - \{0\}) \rightarrow \mathcal{S}(\mathbb{R})$ and $N : \mathcal{S}(\mathbb{R}^{2n}) \rightarrow \mathcal{H}$ are continuous and surjective. Let us consider now $\mu \in \mathcal{S}'(\mathbb{R}^{2n})^{U(p,q)}$, then it is easy to see that there exists a unique $T \in \mathcal{S}'(\mathbb{R})$ such that $\langle \mu, f \rangle = \langle$

$T, Nf >$, for every $f \in \mathcal{S}(\mathbb{R}^{2n} - \{0\})$. Moreover, if $N' : \mathcal{H}' \rightarrow \mathcal{S}'(\mathbb{R}^{2n})$ is the adjoint map, by following again the arguments shown on [T] we can see that N' is a homeomorphism. Finally, for a function $f \in \mathcal{S}(N(p, q, \mathbb{C}))$ we write $Nf(\tau, t)$ for $N(f(., t))(\tau)$.

The distributions $S_{\lambda, k}$ are defined as follows

$$S_{\lambda, k} = \sum_{\{m \in \mathbb{N}_0^n, B(m)=k\}} E_\lambda(h_m, h_m), \quad (1.5)$$

where $B(m) = \sum_{j=1}^p m_j - \sum_{j=p+1}^n m_j$, the set of functions $\{h_m\} \subset L^2(\mathbb{R}^n)$ is the Hermite basis and $E_\lambda(h, h')(z, t) = \langle \pi_\lambda(z, t)h, h' \rangle$ are the matrix entries of the Schrödinger representation π_λ . Also, $S_{\lambda, k} = e^{-i\lambda t} \otimes F_{\lambda, k}$, where each $F_{\lambda, k} \in \mathcal{S}'(\mathbb{C}^n)^{U(p, q)}$ is a tempered distribution defined in terms of the Laguerre polynomials L_k^m and the map N as follows: for $g \in \mathcal{S}(\mathbb{C}^n)$, $\lambda \neq 0$ and $k \in \mathbb{Z}$

$$\text{if } k \geq 0, \langle F_{\lambda, k}, g \rangle = \langle (L_{k-q+n-1}^0 H)^{(n-1)}, \tau \rightarrow \frac{2}{|\lambda|} e^{-\frac{\tau}{2}} N g \left(\frac{2}{|\lambda|} \tau \right) \rangle, \text{ and} \quad (1.6)$$

$$\text{if } k < 0, \langle F_{\lambda, k}, g \rangle = \langle (L_{-k-p+n-1}^0 H)^{(n-1)}, \tau \rightarrow \frac{2}{|\lambda|} e^{-\frac{\tau}{2}} N g \left(-\frac{2}{|\lambda|} \tau \right) \rangle. \quad (1.7)$$

For the quaternionic case we consider the Schrödinger representation π_w as given in [R] (see also [K-R]):

$$\pi_w(\alpha, z) = \pi_{|w|} \left(\alpha, \langle z, \frac{w}{|w|} \rangle \right), \quad (1.8)$$

where $\pi_{|w|}$ is the Schrödinger representation for the classical Heisenberg group $N(2p, 2q, \mathbb{C})$. We have analogously that the distributions $\varphi_{w, k}$ are defined by

$$\varphi_{w, k} = \sum_{m \in \mathbb{N}_0^{2n}: B(m)=k} E_w(h_m, h_m), \quad (1.9)$$

where $B(m) = \sum_{j=1}^{2p} m_j - \sum_{j=2p+1}^{2n} m_j$ and $E_w(h, h')(\alpha, z) = \langle \pi_w(\alpha, z)h, h' \rangle$ are the matrix entries of the Schrödinger representation π_w . Moreover, we have that $\varphi_{w, k} = e^{i\langle w, z \rangle} \otimes \theta_{w, k}$, where $\theta_{w, k}$ is a tempered distribution such that $\theta_{w, k} = N'T_{|w|, k}$. If we set $\lambda = |w|$, we have that $T_{|w|, k} = F_{\lambda, k}$ where we replace n, p, q by $2n, 2p, 2q$, respectively, in (1.6) and (1.7). Observe that if we define

$$\varphi_{\lambda, k}(\alpha, z) = \int_{S^2} e^{i\langle z, \lambda \xi \rangle} d\xi \theta_{\lambda, k}(\alpha), \quad (1.10)$$

this distributions are $Spin(3) \otimes U(p, q, \mathbb{H})$ - invariant.

2 A fundamental solution for the operator \mathcal{L}_α

We have that Φ_α defined as in (1.3) is a well defined tempered distribution, and a fundamental solution for \mathcal{L}_α . We include the proof since a misprint in Lemma 1 of [M-R] is used in the proof Lemma 2.10 of [B-D-R].

We will consider $\alpha \in \mathbb{C}$ such that $2k + p - q \pm \alpha \neq 0$ for all $k \in \mathbb{Z}$.

Theorem 2.1. Φ_α defined as in (1.3) is a well defined tempered distribution and it is a fundamental solution for the operator \mathcal{L}_α .

Proof. From (1.3) and (1.5) we can write

$$\begin{aligned} |<\Phi_\alpha, f>| &\leq \sum_{k \in \mathbb{Z}} \int_0^\infty \left(\left| \frac{< S_{-\lambda, k}, f >}{(2k + p - q + \alpha)} \right| + \left| \frac{< S_{\lambda, k}, f >}{(2k + p - q - \alpha)} \right| \right) |\lambda|^{n-1} d\lambda \leq \\ &\leq \sum_{k \in \mathbb{Z}} \int_0^\infty \sum_{\substack{\beta \in \mathbb{N}_0^n \\ B(\beta)=k}} \left(\left| \frac{< E_{-\lambda}(h_\beta, h_\beta), f >}{(2k + p - q + \alpha)} \right| + \left| \frac{< E_\lambda(h_\beta, h_\beta), f >}{(2k + p - q - \alpha)} \right| \right) |\lambda|^{n-1} d\lambda. \end{aligned}$$

From the known facts that $\sum_{k \in \mathbb{Z}} \sum_{\substack{\beta \in \mathbb{N}_0^n \\ B(\beta)=k}} p(\beta) = \sum_{k \geq 0} \binom{k+n-1}{n-1} p(k)$,

$|< E_\lambda(h_\beta, h_\beta), f >| = |< \pi_\lambda(f)h_\beta, h_\beta >| \leq \|f\|_{L^1(N(p, q, \mathbb{C}))}$ and that for $m \in \mathbb{N}$

$$\pi_\lambda(f)h_\beta = \frac{1}{(-1)^m |\lambda|^m (2B(\beta) + p - q + \alpha \operatorname{sgn}(\lambda))^m} \pi_\lambda(L^m f)h_\beta$$

we get that

$$\begin{aligned} |<\Phi_\alpha, f>| &\leq \|L^m f\|_{L^1(N(p, q, \mathbb{C}))} \times \\ &\times \sum_{k \geq 0} \int_0^\infty \binom{k+n-1}{k} \left(\frac{|\lambda|^{n-1-m}}{|2k + p - q + \alpha|^{m+1}} + \frac{|\lambda|^{n-1-m}}{|2k + p - q - \alpha|^{m+1}} \right) d\lambda. \end{aligned}$$

Let us consider the first term, the second one is analogous. We split the integral between $|\lambda| |2k + p - q + \alpha| \geq 1$ and $0 \leq |\lambda| |2k + p - q + \alpha| \leq 1$. Thus,

$$\sum_{k \geq 0} \binom{k+n-1}{k} \int_{|\lambda| |2k + p - q + \alpha| \geq 1} \frac{1}{|2k + p - q + \alpha|^{m+1}} |\lambda|^{n-1-m} d\lambda$$

is finite if we take $m > n$, and

$$\sum_{k \geq 0} \binom{k+n-1}{k} \int_{0 \leq |\lambda| |2k + p - q + \alpha| \leq 1} \frac{1}{|2k + p - q + \alpha|^{m+1}} |\lambda|^{n-1-m} d\lambda$$

is finite for $m < n$. From the above computations it also follows that Φ_α is a tempered distribution. Next we see that it is a fundamental solution by writing $L = L_0 + L_1$, which in coordinates are

$$\begin{aligned} L_0 &= \frac{1}{4} \left(\sum_{j=1}^p (x_j^2 + y_j^2) - \sum_{j=p+1}^n (x_j^2 + y_j^2) \right) \frac{\partial^2}{\partial t^2} + \\ &+ \sum_{j=1}^p \left(\frac{\partial^2}{\partial x_j^2} + \frac{\partial^2}{\partial y_j^2} \right) - \sum_{j=p+1}^n \left(\frac{\partial^2}{\partial x_j^2} + \frac{\partial^2}{\partial y_j^2} \right), \\ L_1 &= \frac{\partial}{\partial t} \sum_{j=1}^n \left(x_j \frac{\partial}{\partial y_j} - y_j \frac{\partial}{\partial x_j} \right). \end{aligned}$$

Then, as L_0, L_1 and T commute with left translations and also $L_0(g^\vee) = (L_0 g)^\vee$, $L_1(g^\vee) = -(L_1 g)^\vee$ and $T(g^\vee) = -(Tg)^\vee$ we get that

$$(\mathcal{L}f * \Phi_\alpha)(z, t) = <\Phi_\alpha, (L_{(z,t)^{-1}} \mathcal{L}f)^\vee> = <\Phi_\alpha, (L_0 - i\alpha)(L_{(z,t)^{-1}} f)^\vee>,$$

because $L_1\Phi_\alpha = 0$. Hence,

$$\begin{aligned}
(\mathcal{L}_\alpha f * \Phi)(z, t) &= \sum_{k \in \mathbb{Z}_{-\infty}} \int_0^\infty \frac{\langle S_{\lambda, k}, (L_0 - i\alpha T)(L_{(z, t)^{-1}} f)^\vee \rangle}{-|\lambda|(2k + p - q - \alpha \operatorname{sgn} \lambda)} |\lambda|^{n-1} d\lambda = \\
&= \sum_{k \in \mathbb{Z}_{-\infty}} \int_0^\infty \frac{\langle (L_0 + i\alpha T)S_{\lambda, k}, (L_{(z, t)^{-1}} f)^\vee \rangle}{-|\lambda|(2k + p - q - \alpha \operatorname{sgn} \lambda)} |\lambda|^{n-1} d\lambda = \\
&= \sum_{k \in \mathbb{Z}_{-\infty}} \int_0^\infty \langle S_{\lambda, k}, (L_{(z, t)^{-1}} f)^\vee \rangle |\lambda|^{n-1} d\lambda = f(z, t),
\end{aligned}$$

because of the inversion formula. The other one, $f * \mathcal{L}_\alpha(f) = f$, is immediate. \square

Now we proceed with the computation of Φ_α . Given that the series (1.3) defining Φ_α converges absolutely, we can split the sum over $k \in \mathbb{Z}$ into the sums for $k \geq q$, for $k \leq -p$ and for $-p < k < q$. In the first case we change the summation index writing $k = k' + q$ and in the second as well, writing $k = k' - p$. So we get

$$\begin{aligned}
\langle \Phi_\alpha, f \rangle &= (-1) \sum_{k' \geq 0} \frac{1}{2k' + n - \alpha} \int_0^\infty [\langle S_{\lambda, k'+q}, f \rangle - \langle S_{\lambda, -k'-p}, f \rangle] |\lambda|^{n-1} d\lambda + \\
&\quad + (-1) \sum_{k' \geq 0} \frac{1}{2k' + n + \alpha} \int_0^\infty [\langle S_{-\lambda, k'+q}, f \rangle - \langle S_{-\lambda, -k'-p}, f \rangle] |\lambda|^{n-1} d\lambda + \\
&\quad + (-1) \sum_{-p < k < q} \int_0^\infty \left(\frac{\langle S_{-\lambda, k}, f \rangle}{2k + p - q + \alpha} + \frac{\langle S_{\lambda, k}, f \rangle}{2k + p - q - \alpha} \right) |\lambda|^{n-1} d\lambda.
\end{aligned}$$

By Abel's Lemma and the Lebesgue dominated convergence theorem we can write $\Phi_\alpha = \Phi_1 + \Phi_2$ where

$$\begin{aligned}
\langle \Phi_1, f \rangle &= \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} (-1) \sum_{k' \geq 0} \frac{r^{2k'+n-\alpha}}{2k'+n-\alpha} \int_0^\infty e^{-\epsilon|\lambda|} \times \\
&\quad \times [\langle S_{\lambda, k'+q}, f \rangle - \langle S_{\lambda, -k'-p}, f \rangle] |\lambda|^{n-1} d\lambda + \\
&\quad + \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} (-1) \sum_{k' \geq 0} \frac{r^{2k'+n+\alpha}}{2k'+n+\alpha} \int_0^\infty e^{-\epsilon|\lambda|} \times \\
&\quad \times [\langle S_{-\lambda, k'+q}, f \rangle - \langle S_{-\lambda, -k'-p}, f \rangle] |\lambda|^{n-1} d\lambda,
\end{aligned} \tag{2.1}$$

$$\begin{aligned}
\langle \Phi_2, f \rangle &= \lim_{\epsilon \rightarrow 0^+} (-1) \sum_{-p < k < q} \int_0^\infty e^{-\epsilon|\lambda|} \times \\
&\quad \times \left(\frac{\langle S_{-\lambda, k}, f \rangle}{2k + p - q + \alpha} + \frac{\langle S_{\lambda, k}, f \rangle}{2k + p - q - \alpha} \right) |\lambda|^{n-1} d\lambda.
\end{aligned} \tag{2.2}$$

Using that $S_{\lambda,k} = e^{-i\lambda t} \otimes F_{\lambda,k}$ and the computations from [G-S(2)], namely (2.6) to (2.9), we get that

$$\begin{aligned}
\langle \Phi_1, f \rangle &= \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} (-1) \sum_{k \geq 0} \frac{r^{2k+n-\alpha}}{2k+n-\alpha} \int_0^\infty e^{-\epsilon|\lambda|} \int_{-\infty}^\infty e^{-i\lambda t} \times \\
&\quad \times \langle (L_{k+n-1}^0 H)^{(n-1)}, \frac{2}{|\lambda|} e^{-\frac{\tau}{2}} [Nf(\frac{2}{|\lambda|}\tau, t) - Nf(-\frac{2}{|\lambda|}\tau, t)] \rangle dt d\lambda + \\
&+ \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} (-1) \sum_{k \geq 0} \frac{r^{2k+n+\alpha}}{2k+n+\alpha} \int_0^\infty e^{-\epsilon|\lambda|} \int_{-\infty}^\infty e^{i\lambda t} \times \\
&\quad \times \langle (L_{k+n-1}^0 H)^{(n-1)}, \frac{2}{|\lambda|} e^{-\frac{\tau}{2}} [Nf(\frac{2}{|\lambda|}\tau, t) - Nf(-\frac{2}{|\lambda|}\tau, t)] \rangle dt d\lambda,
\end{aligned}$$

and setting

$$b_{k,l} = \sum_{j=l}^{n-2} \binom{j}{l} \left(\frac{1}{2}\right)^{2-l} (-1)^{n-j} \binom{k+n-1}{n-j-2}, \quad (2.3)$$

we have that

$$\begin{aligned}
\langle \Phi_1, f \rangle &= \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} \sum_{k \geq 0} \frac{r^{2k+n-\alpha}}{2k+n-\alpha} \int_0^\infty e^{-\epsilon|\lambda|} \int_{-\infty}^\infty e^{-i\lambda t} \times \\
&\quad \times \left[(-1)^n \int_{-\infty}^\infty L_k^{n-1} \left(\frac{|\lambda|}{2} |s|\right) e^{-\frac{|\lambda|}{4} |s|} \operatorname{sgn}(s) Nf(s, t) ds + \right. \\
&\quad \left. - 2 \sum_{\substack{l=0 \\ l \text{ odd}}}^{n-2} \left(\frac{2}{|\lambda|}\right)^{l+1} b_{k,l} \frac{\partial^l Nf}{\partial \tau^l}(0, t) \right] dt d\lambda + \\
&+ \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} \sum_{k \geq 0} \frac{r^{2k+n+\alpha}}{2k+n+\alpha} \int_0^\infty e^{-\epsilon|\lambda|} \int_{-\infty}^\infty e^{i\lambda t} \times \\
&\quad \times \left[(-1)^n \int_{-\infty}^\infty L_k^{n-1} \left(\frac{|\lambda|}{2} |s|\right) e^{-\frac{|\lambda|}{4} |s|} \operatorname{sgn}(s) Nf(s, t) ds + \right. \\
&\quad \left. - 2 \sum_{\substack{l=0 \\ l \text{ odd}}}^{n-2} \left(\frac{2}{|\lambda|}\right)^{l+1} b_{k,l} \frac{\partial^l Nf}{\partial \tau^l}(0, t) \right] dt d\lambda.
\end{aligned}$$

Now we define

$$G_f(\tau, t) = Nf(\tau, t) - \sum_{j=0}^{n-2} \frac{\partial^j Nf}{\partial \tau^j}(0, t) \frac{\tau^j}{j!}, \quad (2.4)$$

and then we can split $\Phi_1 = \Phi_{11} + \Phi_{12}$ where

$$\begin{aligned}
\langle \Phi_{11}, f \rangle &= \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} \sum_{k \geq 0} (-1)^n \frac{r^{2k+n-\alpha}}{2k+n-\alpha} \int_0^\infty \int_{-\infty}^\infty e^{-\epsilon|\lambda|} e^{-i\lambda t} |\lambda|^{n-1} \times \\
&\quad \times \int_{-\infty}^\infty L_k^{n-1} \left(\frac{|\lambda|}{2} |\tau| \right) e^{-\frac{|\lambda|}{4} |\tau|} \operatorname{sgn}(\tau) G_f(\tau, t) d\tau dt d\lambda + \\
&+ \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} \sum_{k \geq 0} (-1)^n \frac{r^{2k+n+\alpha}}{2k+n+\alpha} \int_0^\infty \int_{-\infty}^\infty e^{-\epsilon|\lambda|} e^{i\lambda t} |\lambda|^{n-1} \times \\
&\quad \times \int_{-\infty}^\infty L_k^{n-1} \left(\frac{|\lambda|}{2} |\tau| \right) e^{-\frac{|\lambda|}{4} |\tau|} \operatorname{sgn}(\tau) G_f(\tau, t) d\tau dt d\lambda,
\end{aligned} \tag{2.5}$$

$$\begin{aligned}
\langle \Phi_{12}, f \rangle &= \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} \sum_{k \geq 0} \frac{r^{2k+n-\alpha}}{2k+n-\alpha} \int_0^\infty \int_{-\infty}^\infty e^{-\epsilon|\lambda|} e^{-i\lambda t} |\lambda|^{n-1} \times \\
&\quad \times 2 \sum_{\substack{l=0 \\ l \text{ odd}}}^{n-2} \left(\frac{2}{|\lambda|} \right)^{l+1} (a_{k,l} + b_{k,l}) \frac{\partial^l N f}{\partial \tau^l}(0, t) dt d\lambda + \\
&+ \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} \sum_{k \geq 0} \frac{r^{2k+n+\alpha}}{2k+n+\alpha} \int_0^\infty \int_{-\infty}^\infty e^{-\epsilon|\lambda|} e^{i\lambda t} |\lambda|^{n-1} \times \\
&\quad \times 2 \sum_{\substack{l=0 \\ l \text{ odd}}}^{n-2} \left(\frac{2}{|\lambda|} \right)^{l+1} (a_{k,l} + b_{k,l}) \frac{\partial^l N f}{\partial \tau^l}(0, t) dt d\lambda,
\end{aligned} \tag{2.6}$$

(2.7)

with

$$a_{k,l} = (-1)^n \frac{1}{l!} \int_0^\infty L_k^{n-1}(s) e^{-\frac{s}{2}} s^l ds. \tag{2.8}$$

We will show that Φ_{11} is well defined. We have proved that the series (1.3) defining Φ_α converges and as Φ_2 is a finite sum we will obtain that Φ_{12} is also well defined.

Proposition 2.2. *The following identities hold:*

(i)

$$\begin{aligned}
&\int_{-\infty}^\infty e^{-\epsilon|\lambda|} e^{-i\lambda t} L_k^{n-1} \left(\frac{|\lambda|}{2} |\tau| \right) e^{-\frac{|\lambda|}{4} |\tau|} |\lambda|^{n-1} d\lambda = \\
&= 4^n (n-1)! (-1)^n \binom{k+n-1}{k} \left(\frac{(|\tau| - 4\epsilon - 4it)^k}{(|\tau| + 4\epsilon + 4it)^{k+n}} \right).
\end{aligned}$$

(ii)

$$\begin{aligned} & \lim_{\epsilon \rightarrow 0^+} \int_{\mathbb{R}^2} \left(\frac{(|\tau| - 4it - 4\epsilon)^k}{(|\tau| + 4it + 4\epsilon)^{k+n}} \right) sgn(\tau) G_f(\tau, t) d\tau dt = \\ &= \int_{\mathbb{R}^2} \frac{1}{(|\tau| - 4it)^{\frac{n}{2} - \frac{\alpha}{2}}} \frac{1}{(|\tau| + 4it)^{\frac{n}{2} + \frac{\alpha}{2}}} \left(\frac{|\tau| - 4it}{\tau^2 + 16t^2} \right)^{2k+n-\alpha} sgn(\tau) G_f(\tau, t) d\tau dt. \end{aligned}$$

(iii)

$$\begin{aligned} & \lim_{\epsilon \rightarrow 0^+} \int_{\mathbb{R}^2} \left(\frac{(|\tau| + 4it - 4\epsilon)^k}{(|\tau| - 4it + 4\epsilon)^{k+n}} \right) sgn(\tau) G_f(\tau, t) d\tau dt = \\ &= \int_{\mathbb{R}^2} \frac{1}{(|\tau| - 4it)^{\frac{n}{2} - \frac{\alpha}{2}}} \frac{1}{(|\tau| + 4it)^{\frac{n}{2} + \frac{\alpha}{2}}} \left(\frac{|\tau| - 4it}{\tau^2 + 16t^2} \right)^{2k+n+\alpha} sgn(\tau) G_f(\tau, t) d\tau dt. \end{aligned}$$

Proof. From (4.9) of [G-S(2)] we know that (i) follows from the generating identity for the Laguerre polynomials:

$$\sum_{k \geq 0} L_k^{n-1}(t) z^k = \frac{1}{(1-z)^n} e^{-\frac{zt}{1-z}}. \quad (2.9)$$

From Lemma 2.2 of [G-S(2)], which states that the function $\frac{G_f(\tau, t)}{(\tau^2 + 16t^2)^{\frac{n}{2}}}$ is integrable in \mathbb{R}^2 and the fact that $\left| \frac{1}{(|\tau| - 4it)^{\frac{n}{2} - \frac{\alpha}{2}}} \right| \left| \frac{1}{(|\tau| + 4it)^{\frac{n}{2} + \frac{\alpha}{2}}} \right| = 1$, it follows that the function $\frac{1}{(|\tau| - 4it)^{\frac{n}{2} - \frac{\alpha}{2}}} \frac{1}{(|\tau| + 4it)^{\frac{n}{2} + \frac{\alpha}{2}}} G_f(\tau, t)$ is integrable in \mathbb{R}^2 . So we get (ii). For (iii) we just change $e^{-i\lambda t}$ by $e^{i\lambda t}$ and argue like for (ii). \square

Then, by Proposition 2.2,

$$\begin{aligned} < \Phi_{11}, f > &= \beta_n \lim_{r \rightarrow 1^-} \sum_{k \geq 0} \frac{r^{2k+n-\alpha}}{2k+n-\alpha} \alpha_k \times \\ & \times \int_{\mathbb{R}^2} \left(\frac{|\tau| - 4it}{\tau^2 + 16t^2} \right)^{2k+n-\alpha} \frac{sgn(\tau) G_f(\tau, t)}{(|\tau| - 4it)^{\frac{n}{2} - \frac{\alpha}{2}} (|\tau| + 4it)^{\frac{n}{2} + \frac{\alpha}{2}}} d\tau dt + \\ & + \beta_n \lim_{r \rightarrow 1^-} \sum_{k \geq 0} \frac{r^{2k+n+\alpha}}{2k+n+\alpha} \alpha_k \times \\ & \times \int_{\mathbb{R}^2} \left(\frac{|\tau| + 4it}{\tau^2 + 16t^2} \right)^{2k+n+\alpha} \frac{sgn(\tau) G_f(\tau, t)}{(|\tau| + 4it)^{\frac{n}{2} + \frac{\alpha}{2}} (|\tau| - 4it)^{\frac{n}{2} - \frac{\alpha}{2}}} d\tau dt, \end{aligned}$$

where $\beta_n = 4^n (n-1)! (-1)^n$ y $\alpha_k = \binom{k+n-1}{k} (-1)^k$.

To study $< \Phi_{11}, f >$ we split each integral over the left and right halfplanes and take polar

coordinates $\tau - 4it = \rho e^{i\theta}$ to obtain

$$\begin{aligned}
\langle \Phi_{11}, f \rangle &= \beta_n \lim_{r \rightarrow 1^-} \sum_{k \geq 0} \alpha_k \frac{r^{2k+n-\alpha}}{2k+n-\alpha} \times \\
&\quad \times \int_0^\infty \left[\int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} e^{i(2k+n-\alpha)\theta} \frac{1}{4\rho^{n-1}} e^{i\alpha\theta} \operatorname{sgn}(\cos \theta) G_f(\rho \cos \theta, -\frac{\rho}{4} \sin \theta) d\theta + \right. \\
&\quad \left. + \int_{\frac{\pi}{2}}^{\frac{3}{2}\pi} \frac{e^{-i(2k+n-\alpha)\theta} e^{-i\alpha\theta}}{(-1)^n 4\rho^{n-1}} \operatorname{sgn}(\cos \theta) G_f(\rho \cos \theta, -\frac{\rho}{4} \sin \theta) d\theta \right] d\rho + \\
&\quad + \beta_n \lim_{r \rightarrow 1^-} \sum_{k \geq 0} \alpha_k \frac{r^{2k+n+\alpha}}{2k+n+\alpha} \times \\
&\quad \times \int_0^\infty \left[\int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} e^{-i(2k+n+\alpha)\theta} \frac{1}{4\rho^{n-1}} e^{i\alpha\theta} \operatorname{sgn}(\cos \theta) G_f(\rho \cos \theta, -\frac{\rho}{4} \sin \theta) d\theta + \right. \\
&\quad \left. + \int_{\frac{\pi}{2}}^{\frac{3}{2}\pi} \frac{e^{i(2k+n+\alpha)\theta} e^{-i\alpha\theta}}{(-1)^n 4\rho^{n-1}} \operatorname{sgn}(\cos \theta) G_f(\rho \cos \theta, -\frac{\rho}{4} \sin \theta) d\theta \right] d\rho
\end{aligned}$$

Now we change variable in the second and fourth term according to $\theta \longleftrightarrow -\theta$. Then, in the fourth term we change variables again according to $\theta \longleftrightarrow \theta + 2\pi$. By proposition 2.2 we can interchange the integration order, so we can write

$$\begin{aligned}
\langle \Phi_{11}, f \rangle &= \beta_n \lim_{r \rightarrow 1^-} \int_0^\infty \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} e^{i\alpha\theta} \times \\
&\quad \times \left[\sum_{k \geq 0} \alpha_k \left(\frac{r^{2k+n-\alpha}}{2k+n-\alpha} e^{i(2k+n-\alpha)\theta} + \frac{r^{2k+n+\alpha}}{2k+n+\alpha} e^{-i(2k+n+\alpha)\theta} \right) \right] \times \\
&\quad \times \frac{1}{\rho^{n-1}} \operatorname{sgn}(\cos \theta) G_f(\rho \cos \theta, -\frac{\rho}{4} \sin \theta) d\theta d\rho + \\
&\quad + \frac{(-1)^n}{4} \beta_n \lim_{r \rightarrow 1^-} \int_0^\infty \int_{\frac{\pi}{2}}^{\frac{3}{2}\pi} e^{i\alpha\theta} \times \\
&\quad \times \left[\sum_{k \geq 0} \alpha_k \left(\frac{r^{2k+n-\alpha}}{2k+n-\alpha} e^{i(2k+n-\alpha)\theta} + \frac{r^{2k+n+\alpha}}{2k+n+\alpha} e^{-i(2k+n+\alpha)\theta} \right) \right] \times \\
&\quad \times \frac{1}{\rho^{n-1}} \operatorname{sgn}(\cos \theta) G_f(\rho \cos \theta, \frac{\rho}{4} \sin \theta) d\theta d\rho.
\end{aligned}$$

Let I denote the real interval $[-\frac{\pi}{2}, \frac{\pi}{2}]$. Next let us consider the vector space

$$\mathcal{X} = \left\{ g \in C^{n-2}(I) : g^{(j)}\left(\pm\frac{\pi}{2}\right) = 0, 0 \leq j \leq n-2, g^{(n-1)} \in L^\infty(I) \right\}.$$

We identify each function $g \in \mathcal{X}$ with the function \tilde{g} on $S^1 = \mathbb{R}/\mathbb{Z}$, defined equal to 0 outside $\text{supp}(g)$ and we make no distinction between g and \tilde{g} . Thus, if $g \in \mathcal{X}$ then $g \in C^{n-2}(S^1)$ with $g^{(n-1)} \in L^\infty(S^1)$. Observe that if $g \in \mathcal{X}$, then also $e^{i\alpha\theta}g \in \mathcal{X}$. The topology on \mathcal{X} is given by $\|g\|_{\mathcal{X}} = \max_{0 \leq j \leq n-1} \|g^{(j)}\|_\infty$.

For $k \in \mathbb{Z}$ we set $\alpha_k = \binom{k+n-1}{k}(-1)^k$. Now let us define

$$\Psi_{r,\alpha}(\theta) = \sum_{k \geq 0} \alpha_k \left(\frac{r^{2k+n-\alpha} e^{i(2k+n-\alpha)\theta}}{2k+n-\alpha} + \frac{r^{2k+n+\alpha} e^{-i(2k+n+\alpha)\theta}}{2k+n+\alpha} \right), \quad (2.10)$$

$$\langle \Psi_\alpha, g \rangle = \langle \sum_{k \geq 0} \alpha_k \left(\frac{e^{i(2k+n-\alpha)\theta}}{2k+n-\alpha} + \frac{e^{-i(2k+n+\alpha)\theta}}{2k+n+\alpha} \right), g \rangle, \quad (2.11)$$

and let us see that $\Psi_\alpha \in \mathcal{X}'$, the dual space of \mathcal{X} . Indeed,

$$|\langle \Psi_\alpha, g \rangle| \leq |e^{i\alpha\theta}| \sum_{k \geq 0} \binom{k+n-1}{k} \left(\frac{|\langle e^{i(2k+n)\theta}, g \rangle|}{|2k+n-\alpha|} + \frac{|\langle e^{-i(2k+n)\theta}, g \rangle|}{|2k+n+\alpha|} \right). \quad (2.12)$$

If $\widehat{g}(n) = \langle g, e^{in\theta} \rangle$ denotes the n -th Fourier coefficient of g , then we have that

$$\begin{aligned} |\langle \Psi_\alpha, g \rangle| &\leq c \sum_{k \geq 0} \frac{k^{n-1}}{|2k+n|^{n-1}} \left(\frac{|\widehat{g}^{(n-1)}(2k+n)|}{|2k+n-\alpha|} + \frac{|\widehat{g}^{(n-1)}(-2k-n)|}{|2k+n+\alpha|} \right) \leq \\ &\leq c \sum_{k \geq 0} \frac{1}{k} |\widehat{g}^{(n-1)}(2k+n)| + \frac{1}{k} |\widehat{g}^{(n-1)}(-2k-n)| \leq \\ &\leq c \left(\sum_{k \geq 0} \frac{1}{k^2} \right)^{\frac{1}{2}} \|\widehat{g}^{(n-1)}\|_{L^2}, \end{aligned}$$

where we used the Cauchy-Schwarz inequality. Observe that the constants c are not the same on each expression.

By Abel's Lemma, $\lim_{r \rightarrow 1^-} \Psi_{r,\alpha} = \Psi_\alpha$ in \mathcal{X}' that is, with respect to the weak convergence topology. Similarly, if J denotes the real interval $[\frac{\pi}{2}, \frac{3}{2}\pi]$, we define the space,

$$\mathcal{Y} = \left\{ g \in C^{n-2}(J) : g^{(j)}\left(\frac{\pi}{2}\right) = g^{(j)}\left(\frac{3}{2}\pi\right) = 0, 0 \leq j \leq n-2, g^{(n-1)} \in L^\infty(J) \right\},$$

and we obtain that Ψ_α is well defined in \mathcal{Y}' and $\lim_{r \rightarrow 1^-} \Psi_{r,\alpha} = \Psi_\alpha$ in \mathcal{Y}' .

Our aim now is to compute Ψ_α .

From Proposition 3.7 of [G-S(2)] we know that if $\Theta \in \mathcal{D}'(S^1)$ is defined by

$$\Theta(\theta) = i \sum_{k \geq 0} \binom{k+n-1}{k} (-1)^k e^{i(2k+n)\theta}, \quad (2.13)$$

then for even n we have that

$$\Re \Theta(\theta) = \frac{d}{d\theta} Q_{n-2} \left(\frac{d}{d\theta} \right) (\delta_{\frac{\pi}{2}} + \delta_{-\frac{\pi}{2}}) = \sum_{j=0}^{n-2} c_j \left(\delta_{\frac{\pi}{2}}^{(j+1)} + \delta_{-\frac{\pi}{2}}^{(j+1)} \right), \quad (2.14)$$

where Q_{n-2} is a polynomial of degree $n-2$; and for odd n we have that

$$\Re \Theta(\theta) = d_0 \frac{d}{d\theta} \tilde{H} + \frac{d}{d\theta} Q_{n-2} \left(\frac{d}{d\theta} \right) (\delta_{\frac{\pi}{2}} - \delta_{-\frac{\pi}{2}}) = d_0 (\delta_{-\frac{\pi}{2}} - \delta_{\frac{\pi}{2}}) + \sum_{j=0}^{n-2} c_j (\delta_{\frac{\pi}{2}}^{(j+1)} - \delta_{-\frac{\pi}{2}}^{(j+1)}), \quad (2.15)$$

where Q_{n-2} is a polynomial of degree $n-2$, and $\tilde{H}(\theta) = H(\cos \theta)$.

Let us recall the generating identity for the Laguerre polynomials (2.9), and take $t = 0$ and $z = -r^2 e^{2i\theta}$. We get

$$\sum_{k \geq 0} \binom{k+n-1}{k} (-1)^k r^{2k+n} e^{i(2k+n)\theta} = \left(\frac{re^{i\theta}}{1+r^2 e^{2i\theta}} \right)^n. \quad (2.16)$$

We also need a couple of results:

Lemma 2.3. *For a fixed $r > 1$ the functions $\alpha \rightarrow \Psi_{r,\alpha}(0)$ and $\alpha \rightarrow \lim_{r \rightarrow 1^-} \Psi_{r,\alpha}(0)$ are analytic on $\Omega = \mathbb{C} \setminus F$, where $F = \{2k+n : k \in \mathbb{Z}\}$.*

Proof. Let K be a compact set, $K \subset \Omega$. It is easy to see that for fixed r the series (2.10) converges uniformly, since

$$|\Psi_{r,\alpha}(0)| \leq \max_{\alpha \in K} |r^\alpha| \left(\frac{r}{1+r^2} \right)^n d(K, F).$$

Also, for $\alpha \in \Omega$ there exists $\lim_{r \rightarrow 1^-} \Psi_{r,\alpha}(0)$. Indeed, if $0 \leq r_1 < r < r_2 < 1$, from the Mean Value Theorem we have that for some $\xi \in (r_1, r_2)$,

$$\begin{aligned} \Psi_{r_1,\alpha}(0) - \Psi_{r_2,\alpha}(0) &= \frac{d}{dr} \Psi_{\xi,\alpha}(0)(r_2 - r_1) = (\xi^{-\alpha-1} + \xi^{\alpha-1}) \sum_{k \geq 0} \alpha_k \xi^{2k+n} (r_2 - r_1) = \\ &= (\xi^{-\alpha-1} + \xi^{\alpha-1}) \left(\frac{\xi}{1+\xi^2} \right)^n (r_2 - r_1), \end{aligned}$$

where the last equality holds from (2.16). Hence

$$|\Psi_{r_1,\alpha}(0) - \Psi_{r_2,\alpha}(0)| \leq c(\xi) |r_2 - r_1|,$$

where $c(\xi)$ is a constant which depends on ξ . Moreover, for $\alpha \in K$ and $\xi \in [\frac{1}{2}, 1]$, $\xi^{n-\alpha-1} + \xi^{n+\alpha-1}$ is bounded in $K \times [\frac{1}{2}, 1]$, so the convergence is uniform, hence $\alpha \rightarrow \lim_{r \rightarrow 1^-} \Psi_{r,\alpha}(0)$ is an analytic function. \square

Lemma 2.4. *Let $0 < \delta < \frac{\pi}{4}$. For $0 < r < 1$ and $0 \leq |\theta| < \delta$ we have that*

$$|\Psi_{r,\alpha}(\theta) - \Psi_{r,\alpha}(0)| \leq \left(\max_{0 \leq |\theta| < \delta} e^{|\Im \alpha| |\theta|} \right) (a|r^{-\alpha} - r^\alpha| + b|r^\alpha|(1-r)) |\theta|,$$

with a, b positive constants. Also for $0 \leq |\theta - \pi| < \delta < \frac{\pi}{4}$,

$$|\Psi_{r,\alpha}(\theta) - \Psi_{r,\alpha}(\pi)| \leq \left(\max_{0 \leq |\theta - \pi| < \delta} e^{|\Im \alpha| |\theta|} \right) (a|r^{-\alpha} - r^\alpha| + b|r^\alpha|(1-r)) |\theta - \pi|,$$

with a, b positive constants.

Proof. We will estimate $|\Psi_{r,\alpha}(\theta) - \Psi_{r,\alpha}(0)|$ for $0 < |\theta| < \delta < \frac{\pi}{4}$, the other case is similar. We have

$$\begin{aligned} \frac{d}{d\theta} \Psi_{r,\alpha}(\theta) &= ie^{-i\alpha\theta} \sum_{k \geq 0} \alpha_k r^{2k+n} \left((r^{-\alpha} - r^\alpha) e^{i(2k+n)\theta} + (e^{i(2k+n)\theta} - e^{-i(2k+n)\theta}) r^\alpha \right) = \\ &= ie^{-i\alpha\theta} \left((r^{-\alpha} - r^\alpha) \left(\frac{re^{i\theta}}{1+r^2 e^{2i\theta}} \right)^n + 2ir^\alpha \Im \left(\frac{re^{i\theta}}{1+r^2 e^{2i\theta}} \right)^n \right), \end{aligned}$$

because of (2.16). We have that

$$\left| \frac{d}{d\theta} \Psi_{r,\alpha}(\theta) \right| \leq e^{|\Im \alpha| |\theta|} \left(|r^{-\alpha} - r^\alpha| \left| \left(\frac{re^{i\theta}}{1+r^2 e^{i2\theta}} \right)^n \right| + 2|r^\alpha| \left| \Im \left(\frac{re^{i\theta}}{1+r^2 e^{i2\theta}} \right)^n \right| \right). \quad (2.17)$$

From Proposition 3.1 of [G-S(2)] we know that $\left| \Im \left(\frac{re^{i\theta}}{1+r^2 e^{i2\theta}} \right)^n \right| \rightarrow 0$ as $r \rightarrow 1^-$, uniformly for $|\theta| < \frac{\pi}{4}$, $|\theta - \pi| < \frac{\pi}{4}$. Also, $\left| \left(\frac{re^{i\theta}}{1+r^2 e^{i2\theta}} \right)^n \right| \leq c$, for a constant c . Then $\left| \frac{d}{d\theta} \Psi_{r,\alpha}(\theta) \right| \rightarrow 0$ uniformly on $|\theta| < \frac{\pi}{4}$ as $r \rightarrow 1^-$, and we get the desired inequality by applying the Mean Value Theorem around 0. \square

Now we can state the following

Proposition 2.5. For $f \in \mathcal{X}$ we have that $\langle \Psi_\alpha, f \rangle = C_\alpha \langle 1, f \rangle$, where $C_\alpha = \frac{\Gamma(\frac{n+\alpha}{2})\Gamma(\frac{n-\alpha}{2})}{(n-1)!}$; and for $f \in \mathcal{Y}$ we have that $\langle \Psi_\alpha, f \rangle = \widetilde{C}_\alpha \langle 1, f \rangle$, where $\widetilde{C}_\alpha = (-1)^n e^{-i\alpha\pi} C_\alpha$.

Proof. First we consider $f \in \mathcal{X}$ such that $\int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} f(t) dt = 0$ and we define $F(\theta) = \int_{-\frac{\pi}{2}}^{\theta} f(t) dt$. It is easy to see that $F \in \mathcal{X}$ and $F' = f$. Because of the integration by parts formula we have that

$$\begin{aligned} \langle \Psi_\alpha, f \rangle &= \langle \Psi_\alpha, F' \rangle = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} \sum_{k \geq 0} \alpha_k \left(\frac{e^{i(2k+n-\alpha)\theta}}{2k+n-\alpha} + \frac{e^{-i(2k+n+\alpha)\theta}}{2k+n+\alpha} \right) F'(\theta) d\theta = \\ &= -\langle \Theta, e^{-i\alpha\theta} F \rangle - \langle \overline{\Theta}, e^{-i\alpha\theta} F \rangle, \end{aligned}$$

where $\overline{\Theta} = \sum_{k \geq 0} \binom{k+n-1}{k} (-1)^k e^{-i(2k+n)\theta}$. So if n is even, from (2.14) we get that

$$\langle \Psi_\alpha, f \rangle = - \sum_{j=0}^{n-2} c_j \langle \delta_{\frac{\pi}{2}}^{(j+1)} + \delta_{-\frac{\pi}{2}}^{(j+1)}, e^{-i\alpha\theta} F \rangle - \sum_{j=0}^{n-2} \overline{c_j} \langle \overline{\delta_{\frac{\pi}{2}}^{(j+1)}} + \overline{\delta_{-\frac{\pi}{2}}^{(j+1)}}, e^{-i\alpha\theta} F \rangle,$$

and because $\langle \delta_{\pm\frac{\pi}{2}}^{(j+1)}, e^{-i\alpha\theta} F \rangle = 0$ we conclude that $\langle \Psi_\alpha, f \rangle = 0$. If n is odd we use (2.15) to conclude that $\langle \Psi_\alpha, f \rangle = 0$.

For the general case of any $f \in \mathcal{X}$ we consider $h \in \mathcal{X}$ such that $\int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} h(t) dt = 1$ and define $g(\theta) = f(\theta) - \left(\int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} f(t) dt \right) h(\theta)$. So we can apply the above result to g and get that $\langle \Psi_\alpha, g \rangle = 0$. Then $\langle \Psi_\alpha, f \rangle = \langle \Psi_\alpha, g \rangle + \langle \Psi_\alpha, h \rangle \langle 1, f \rangle = \langle \Psi_\alpha, h \rangle \langle 1, f \rangle$. Let $C_\alpha = \langle \Psi_\alpha, h \rangle$.

In order to compute C_α , consider $g \in \mathcal{X}$ such that $\text{supp}(g) \subset (-\frac{\pi}{4}, \frac{\pi}{4})$, $\int_{-\frac{\pi}{4}}^{\frac{\pi}{4}} g(t) dt = 1$ and $g \geq 0$

we have that

$$\langle e^{i\alpha\theta} \Psi_\alpha, g \rangle = C_\alpha \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} e^{i\alpha\theta} g(\theta) d\theta,$$

and also that

$$\langle e^{i\alpha\theta} \Psi_\alpha, g \rangle = \lim_{r \rightarrow 1^-} \left(\int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} (\Psi_{r,\alpha}(\theta) - \Psi_{r,\alpha}(0)) e^{i\alpha\theta} g(\theta) d\theta + \Psi_{r,\alpha}(0) \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} e^{i\alpha\theta} g(\theta) d\theta \right).$$

From lemmas 1 and 2 we get that

$$C_\alpha = \lim_{r \rightarrow 1^-} \Psi_{r,\alpha}(0)$$

and also that C_α is an analytic function of α . Given that $\Psi_{0,\alpha}(0) = 0$ we can write

$$C_\alpha = \lim_{r \rightarrow 1^-} \Psi_{r,\alpha}(0) = \Psi_{1,\alpha}(0) - \Psi_{0,\alpha}(0) = \int_0^1 w'_\alpha(s) ds,$$

where

$$w_\alpha(r) = \Psi_{r,\alpha}(0) = r^{-\alpha} \sum_{k \geq 0} \alpha_k \frac{r^{2k+n}}{2k+n-\alpha} + r^\alpha \sum_{k \geq 0} \alpha_k \frac{r^{2k+n}}{2k+n+\alpha}.$$

Applying (2.9) with $\theta = 0$ we obtain $w'_\alpha(r) = (r^{-\alpha-1} + r^{\alpha-1}) \sum_{k \geq 0} \alpha_k r^{2k+n} = (r^{-\alpha-1} + r^{\alpha-1}) \left(\frac{r}{1+r^2} \right)^n$,

and we can solve the integral for $\Re(n+\alpha) > 0$, $\Re(n-\alpha) > 0$, getting

$$C_\alpha = B \left(\frac{n+\alpha}{2}, \frac{n-\alpha}{2} \right) = \frac{\Gamma \left(\frac{n+\alpha}{2} \right) \Gamma \left(\frac{n-\alpha}{2} \right)}{(n-1)!}, \quad (2.18)$$

where B is the *Beta function* and Γ is the *Gamma function*. By lemma 2.3, (2.18) holds for the other range of α , by analytic continuation. In a completely analogous way we get that $\widetilde{C}_\alpha = (-1)^n e^{-i\alpha\pi} C_\alpha$. \square

Let us now define

$$K_{1f}(\rho, \theta) = \frac{1}{\rho^{n-1}} \operatorname{sgn}(\cos \theta) G_f(\rho \cos \theta, -\frac{\rho}{4} \sin \theta), \quad (2.19)$$

for $\theta \in [-\frac{\pi}{2}, \frac{\pi}{2}]$, $0 < \rho < \infty$, where G_f is the function defined in (2.4); and

$$K_{2f}(\rho, \theta) = \frac{1}{\rho^{n-1}} \operatorname{sgn}(\cos \theta) G_f(\rho \cos \theta, \frac{\rho}{4} \sin \theta), \quad (2.20)$$

for $\theta \in [\frac{\pi}{2}, \frac{3\pi}{2}]$, $0 < \rho < \infty$.

It is easy to check that $K_{1f}(\rho, \cdot) \in \mathcal{X}$. Recall that we changed variables according to $\tau - 4it = \rho e^{i\theta}$. Since $Nf \in \mathcal{H}_n$, there exists a positive constant c such that

$$\sup_{\tau \neq 0, t \in \mathbb{R}} |(\tau^2 + 16t^2)Nf(\tau, t)| \leq c,$$

that is

$$\left| Nf(\rho \cos \theta, -\frac{\rho}{4} \sin \theta) \right| \leq \frac{c}{\rho^2}.$$

Also since $Nf(0, \cdot) \in \mathcal{S}(\mathbb{R})$, there exists a positive constant c_N such that for $t \in \mathbb{R}$,

$$\left| t^N \sum_{j=0}^{n-2} \frac{\partial^j}{\partial \tau^j} Nf(0, t) \frac{\tau^j}{j!} \right| \leq c_N |\tau|^{n-2}.$$

Thus, for $N \in \mathbb{N}$ there exists c_N such that

$$|K_{1f}(\rho, \theta)| \leq \frac{a}{\rho^{n+1}} + \frac{b}{\rho^{N+1}} \frac{|\cos \theta|^{n-2}}{|\sin \theta|^N}. \quad (2.21)$$

Analogous observations are also true for K_{2f} .

Proposition 2.6. Let C_α and \widetilde{C}_α be the constants obtained in (2.18). Let K_{1f} and K_{2f} be as defined in (2.19) and (2.20) and $\alpha_k = \binom{k+n-1}{k}(-1)^k$. Then

$$\begin{aligned} & \lim_{r \rightarrow 1^-} \int_0^\infty \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} e^{i\alpha\theta} \sum_{k \geq 0} \alpha_k \left(\frac{r^{2k+n-\alpha} e^{i(2k+n-\alpha)\theta}}{2k+n-\alpha} + \frac{r^{2k+n+\alpha} e^{-i(2k+n+\alpha)\theta}}{2k+n+\alpha} \right) K_{1f}(\rho, \theta) d\theta d\rho = \\ & = 4^{n-1} (n-1)! C_\alpha \int_{\mathbb{R}} \int_{\tau > 0} \frac{1}{(\tau - 4it)^{\frac{n-\alpha}{2}}} \frac{1}{(\tau + 4it)^{\frac{n+\alpha}{2}}} \operatorname{sgn}(\tau) G_f(\tau, t) d\tau dt, \end{aligned}$$

and

$$\begin{aligned} & \lim_{r \rightarrow 1^-} \int_0^{\frac{3\pi}{2}} \int_{\frac{\pi}{2}}^{\frac{3\pi}{2}} e^{i\alpha\theta} \sum_{k \geq 0} \alpha_k \left(\frac{r^{2k+n-\alpha} e^{i(2k+n-\alpha)\theta}}{2k+n-\alpha} + \frac{r^{2k+n+\alpha} e^{-i(2k+n+\alpha)\theta}}{2k+n+\alpha} \right) K_{2f}(\rho, \theta) d\theta d\rho = \\ & = 4^{n-1} (n-1)! \widetilde{C}_\alpha \int_{\mathbb{R}} \int_{\tau < 0} \frac{1}{(\tau - 4it)^{\frac{n-\alpha}{2}}} \frac{1}{(\tau + 4it)^{\frac{n+\alpha}{2}}} \operatorname{sgn}(\tau) G_f(\tau, t) d\tau dt. \end{aligned}$$

Proof. The proof follows the same lines that Proposition 4.2 of [G-S(2)]. We only sketch it for completeness.

Taking polar coordinates $\tau - 4it = \rho e^{i\theta}$ we only need to show that

$$\lim_{r \rightarrow 1^-} \int_0^\infty \langle \Psi_{r,\alpha}, e^{i\alpha\theta} K_{1f}(\rho, \theta) \rangle d\rho = \int_0^\infty \langle C_\alpha, e^{i\alpha\theta} K_{1f}(\rho, \theta) \rangle d\rho. \quad (2.22)$$

In order to do this we split the integral for $0 < \rho < 1$ and $1 < \rho < \infty$.

We consider first the case $1 < \rho < \infty$. For $|\theta| \leq \delta < \frac{\pi}{4}$, set $I = \int_1^\infty \int_{|\theta| < \delta} e^{i\alpha\theta} (\Psi_{r,\alpha}(\theta) - \Psi_{r,\alpha}(0)) K_{1f}(\rho, \theta) d\theta d\rho$ and $II = \int_1^\infty \int_{|\theta| < \delta} e^{i\alpha\theta} (\Psi_{r,\alpha}(0) - C_\alpha) K_{1f}(\rho, \theta) d\theta d\rho$. We bound I close to 0 by applying Lemma 2.4 and taking $N = 1$ in (2.21). For II we just take $N = \frac{1}{2}$ in (2.21). To analyze the case $\delta \leq |\theta| \leq \frac{\pi}{2}$, we observe that the function $K_{1f}^*(\theta) = \int_1^\infty K_{1f}(\rho, \theta) d\rho$ defined for $\theta \in [-\frac{\pi}{2}, -\delta] \cup [\delta, \frac{\pi}{2}]$ can be extended to an element of \mathcal{X} that we still denote by K_{1f}^* . Then

$$\begin{aligned} & \int_1^\infty \int_{\delta < |\theta| < \frac{\pi}{2}} e^{i\alpha\theta} (\Psi_{r,\alpha}(\theta) - C_\alpha) K_{1f}(\rho, \theta) d\theta d\rho = \\ & = \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} e^{i\alpha\theta} (\Psi_{r,\alpha}(\theta) - C_\alpha) K_{1f}^*(\theta) d\theta - \int_{|\theta| < \delta} e^{i\alpha\theta} (\Psi_{r,\alpha}(\theta) - C_\alpha) K_{1f}^*(\theta) d\theta. \end{aligned}$$

The first term converges to zero as $r \rightarrow 1^-$ since $\Psi_{r,\alpha} \rightarrow C_\alpha$ as $r \rightarrow 1^-$ in \mathcal{X}' . For the second term we argue as above.

Finally, for the case $0 < \rho < 1$ we apply the same arguments to the function $K_{1f}^{**}(\theta) = \int_0^1 K_{1f}(\rho, \theta) d\rho$. □

Corollary 2.7. $\langle \Phi_{11}, f \rangle$ is well defined for $f \in \mathcal{S}(\mathbb{H}_n)$, and

$$\begin{aligned} \langle \Phi_{11}, f \rangle &= 4^{n-1}(n-1)!C_\alpha \int_{\mathbb{R}} \int_{\tau>0} \frac{1}{(\tau-4it)^{\frac{n-\alpha}{2}}} \frac{1}{(\tau+4it)^{\frac{n+\alpha}{2}}} \operatorname{sgn}(\tau) G_f(\tau, t) d\tau dt + \\ &+ 4^{n-1}(n-1)!\widetilde{C}_\alpha \int_{\mathbb{R}} \int_{\tau<0} \frac{1}{(\tau-4it)^{\frac{n-\alpha}{2}}} \frac{1}{(\tau+4it)^{\frac{n+\alpha}{2}}} \operatorname{sgn}(\tau) G_f(\tau, t) d\tau dt. \end{aligned}$$

From the corollary we also get that $\langle \Phi_{12}, f \rangle$ is well defined. In order to explicitly compute it we define for $0 \leq l \leq n-2$, $\epsilon > 0$ and $f \in \mathcal{S}(\mathbb{H}_n)$

$$d_{\epsilon, l, f}^- = \int_0^\infty \int_{-\infty}^\infty e^{-\epsilon|\lambda|} e^{-i\lambda t} |\lambda|^{n-l-2} \frac{\partial^l}{\partial \tau^l} Nf(0, t) dt d\lambda, \quad \text{and} \quad (2.23)$$

$$d_{\epsilon, l, f}^+ = \int_0^\infty \int_{-\infty}^\infty e^{-\epsilon|\lambda|} e^{i\lambda t} |\lambda|^{n-l-2} \frac{\partial^l}{\partial \tau^l} Nf(0, t) dt d\lambda. \quad (2.24)$$

Then we can write (2.7) as

$$\langle \Phi_{12}, f \rangle = \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} \sum_{k \geq 0} \sum_{\substack{l=0 \\ l \text{ odd}}}^{n-2} 2^{l+2} (a_{kl} + b_{kl}) \left[\frac{r^{2k+n-\alpha}}{2k+n-\alpha} d_{\epsilon, l, f}^- + \frac{r^{2k+n+\alpha}}{2k+n+\alpha} d_{\epsilon, l, f}^+ \right].$$

From Lemma 4.4 in [G-S(2)] we deduce that

$$a_{kl} + b_{kl} = (-1)^k \sum_{j=1}^{l+1} \frac{1}{2^{n-l-j-1}} \binom{n-j-1}{l-j+1} \binom{j+k-1}{k}.$$

Also we have the following

Lemma 2.8. If $0 \leq l \leq n-2$, $\epsilon > 0$ and $f \in \mathcal{S}(\mathbb{H}_n)$, then

$$\lim_{\epsilon \rightarrow 0^+} d_{\epsilon, l, f}^- = \frac{1}{i^{n-l-2}} \langle \frac{\pi}{2} \delta - i(vp) \left(\frac{1}{\lambda} \right), \frac{\partial^{n-2}}{\partial \lambda^{n-l-2} \partial \tau^l} Nf(0, \cdot) \rangle \quad \text{and}$$

$$\lim_{\epsilon \rightarrow 0^+} d_{\epsilon, l, f}^+ = i^{n-l-2} \langle \frac{\pi}{2} \delta + i(vp) \left(\frac{1}{\lambda} \right), \frac{\partial^{n-2}}{\partial \lambda^{n-l-2} \partial \tau^l} Nf(0, \cdot) \rangle.$$

Proof. Let us consider $g(\lambda) = e^{-\epsilon|\lambda|} |\lambda|^{n-l-2}$ and $h(t) = \frac{\partial^l}{\partial \tau^l} Nf(0, t)$, and observe that $\int_{-\infty}^\infty e^{-i\lambda t} h(t) dt = \hat{h}(\lambda)$. Then just by using properties of the Fourier transform we get that

$$d_{\epsilon, l, f}^- = \int_0^\infty \int_{-\infty}^\infty g(\lambda) e^{-i\lambda t} h(t) dt d\lambda = \int_0^\infty g(\lambda) \hat{h}(\lambda) d\lambda = \frac{1}{i^{n-l-2}} \int_{-\infty}^\infty \frac{1}{\epsilon + i\lambda} h^{(n-l-2)}(\lambda) d\lambda.$$

For each $\epsilon > 0$, $\frac{1}{\epsilon + i\lambda}$ is a distribution such that there exists $\lim_{\epsilon \rightarrow 0^+} \frac{1}{\epsilon + i\lambda}$ in $\mathcal{S}'(\mathbb{R})$. Moreover, it is easy to check that

$$\lim_{\epsilon \rightarrow 0^+} \frac{1}{\epsilon + i\lambda} = \frac{\pi}{2} \delta - i(vp) \left(\frac{1}{\lambda} \right).$$

Thus the desired equality follows. For $d_{\epsilon, l, f}^+$ we need to change variable according to $\lambda \longleftrightarrow -\lambda$ after considering the Fourier transform of h . \square

We define for $j \in \mathbb{N}$, $0 < j < n - 1$, the functions of r , with $0 \leq r < 1$, by

$$w_j^-(r) = \sum_{k \geq 0} (-1)^k \binom{j+k-1}{k} \frac{r^{2k+n-\alpha}}{2k+n-\alpha},$$

$$w_j^+(r) = \sum_{k \geq 0} (-1)^k \binom{j+k-1}{k} \frac{r^{2k+n+\alpha}}{2k+n+\alpha}.$$

We can see, in a complete analogous way as the computations made for C_α and \widetilde{C}_α , that this functions are well defined and that

$$c_j^- := \lim_{r \rightarrow 1^-} w_j^-(r) = \frac{1}{2} B_{\frac{1}{2}} \left(\frac{n-\alpha}{2}, j - \frac{n-\alpha}{2} \right), \text{ and}$$

$$c_j^+ := \lim_{r \rightarrow 1^-} w_j^+(r) = \frac{1}{2} B_{\frac{1}{2}} \left(\frac{n+\alpha}{2}, j - \frac{n+\alpha}{2} \right), \quad (2.25)$$

where $B_{\frac{1}{2}}$ is another special function called the *incomplete Beta function*.

We now plug all of this definitions and results together to finally obtain an expression for Φ_{12} :

$$\begin{aligned} <\Phi_{12}, f> &= \sum_{\substack{l=0 \\ l \text{ odd}}}^{n-2} \sum_{j=1}^{l+1} 2^{2l-n+j+3} \binom{n-j-1}{l-j+1} \left[\left(\frac{1}{i^{n-l-2}} c_j^- + i^{n-l-2} c_j^+ \right) \frac{\pi}{2} \right] \times \\ &\quad \times <\delta, \frac{\partial^{n-2}}{\partial \lambda^{n-l-2} \partial \tau^l} Nf(0,.)> + \\ &+ (-1) \sum_{\substack{l=0 \\ l \text{ odd}}}^{n-2} \sum_{j=1}^{l+1} 2^{2l-n+j+3} \binom{n-j-1}{l-j+1} \left(\frac{1}{i^{n-l+1}} c_j^- + i^{n-l+1} c_j^+ \right) \times \\ &\quad \times <(vp) \left(\frac{1}{\lambda} \right), \frac{\partial^{n-2}}{\partial \lambda^{n-l-2} \partial \tau^l} Nf(0,.)>. \end{aligned}$$

All we need to do now is to use again lemma 2.8 to get an expression for Φ_2 . Thus, we have proved the following

Theorem 2.9. *Let C_α and \widetilde{C}_α be the constans defined as in (2.18), and let c_j^- and c_j^+ the constants*

defined as in (2.25). Then

$$\begin{aligned}
\langle \Phi, f \rangle &= 4^{n-1}(n-1)!C_\alpha \int_{\mathbb{R}} \int_{\tau>0} \frac{1}{(\tau-4it)^{\frac{n-\alpha}{2}}} \frac{1}{(\tau+4it)^{\frac{n+\alpha}{2}}} \operatorname{sgn}(\tau) G_f(\tau, t) d\tau dt + \\
&+ 4^{n-1}(n-1)!\widetilde{C}_\alpha \int_{\mathbb{R}} \int_{\tau<0} \frac{1}{(\tau-4it)^{\frac{n-\alpha}{2}}} \frac{1}{(\tau+4it)^{\frac{n+\alpha}{2}}} \operatorname{sgn}(\tau) G_f(\tau, t) d\tau dt + \\
&+ \sum_{\substack{l=0 \\ l \text{ odd}}}^{n-2} \sum_{j=1}^{l+1} 2^{2l-n+j+3} \binom{n-j-1}{l-j+1} \left[\left(\frac{1}{i^{n-l-2}} c_j^- + i^{n-l-2} c_j^+ \right) \frac{\pi}{2} \right] \times \\
&\times \langle \delta, \frac{\partial^{n-2}}{\partial \lambda^{n-l-2} \partial \tau^l} Nf(0, \cdot) \rangle + \\
&+ (-1) \sum_{\substack{l=0 \\ l \text{ odd}}}^{n-2} \sum_{j=1}^{l+1} 2^{2l-n+j+3} \binom{n-j-1}{l-j+1} \left(\frac{1}{i^{n-l+1}} c_j^- + i^{n-l+1} c_j^+ \right) \times \\
&\times \langle (vp) \left(\frac{1}{\lambda} \right), \frac{\partial^{n-2}}{\partial \lambda^{n-l-2} \partial \tau^l} Nf(0, \cdot) \rangle + \\
&+ \sum_{k=0}^{n-2} \sum_{l=0}^{n-2} c_{kl} \left[\frac{1}{n-2k+\alpha-2} i^{n-l-2} \langle \frac{\pi}{2} \delta + i(vp) \left(\frac{1}{\lambda} \right), \frac{\partial^{n-2}}{\partial \lambda^{n-l-2} \partial \tau^l} Nf(0, \cdot) \rangle + \right. \\
&\left. + \frac{1}{n-2k-\alpha-2} \frac{1}{i^{n-l-2}} \langle \frac{\pi}{2} \delta - i(vp) \left(\frac{1}{\lambda} \right), \frac{\partial^{n-2}}{\partial \lambda^{n-l-2} \partial \tau^l} Nf(0, \cdot) \rangle \right],
\end{aligned}$$

where $c_{kl} = \sum_{\substack{1 \leq j \leq n-2 \\ j \geq n-k-2}} 2^{2l-n-j} (-1)^{n-j} \binom{j}{l} \binom{k-l+1}{n-j-2}$.

3 A fundamental solution for L

Like in the classical case, we have that the distribution Φ defined in (1.4) is a well defined tempered distribution and it is a relative fundamental solution for the operator L . The proof is identical to the one of theorem 2.1.

We will compute the fundamental solution Φ by means of the Radon transform and the fundamental solution of the operator L in the classical case \mathbb{H}_{2n} .

Let $F \in \mathcal{S}(\mathbb{R}^3)$. We assign to F a function $\mathcal{R}F : \mathbb{R} \times S^2 \rightarrow \mathbb{R}$ defined by

$$\mathcal{R}F(t, \xi) = \int_{\mathbb{R}^2} F(t\xi + u_1 e_1 + u_2 e_2) du_1 du_2,$$

where $\{\xi, e_1, e_2\}$ is an orthonormal basis of \mathbb{R}^3 . It is easy to see that this definition does not depend on the choice of the basis. In order to recover F from $\mathcal{R}F$, we consider the space $\mathcal{S}(\mathbb{R} \times S^2)$ of the continuous functions $G : \mathbb{R} \times S^2 \rightarrow \mathbb{R}$ that are infinitely differentiable in t and satisfy for every $m, k \in \mathbb{N}_0$ that

$$\sup_{t \in \mathbb{R}, \xi \in S^2} \left| t^m \frac{\partial^k}{\partial t^k} G(t, \xi) \right| < \infty.$$

Now for $G \in \mathcal{S}(\mathbb{R} \times S^2)$ we define a function $\mathcal{R}^*G : \mathbb{R}^3 \rightarrow \mathbb{R}$ by

$$\mathcal{R}^*G(z) = \int_{S^2} G(\langle z, \xi \rangle, \xi) d\xi.$$

Both assignments are well defined, and moreover $\mathcal{R} : \mathcal{S}(\mathbb{R}^3) \rightarrow \mathcal{S}(\mathbb{R} \times S^2)$ is the *Radon transform*, $\mathcal{R}^* : \mathcal{S}(\mathbb{R} \times S^2) \rightarrow \mathcal{S}(\mathbb{R}^3)$ is the *dual Radon transform* and they satisfy for every $F \in \mathcal{S}(\mathbb{R}^3)$

$$-2\pi F = \Delta \mathcal{R}^* \mathcal{R} F, \quad (3.1)$$

where $\Delta = \frac{\partial^2}{\partial z_1^2} + \frac{\partial^2}{\partial z_2^2} + \frac{\partial^2}{\partial z_3^2}$ is the \mathbb{R}^3 -Laplacian.

Now let us consider the function ϕ defined for a fixed $\tau, \tau \neq 0$ by

$$\phi(\tau, z) = \frac{16n}{\pi} \frac{4^{2n}(2n-1)!c_0}{(\tau^2 + 16|z|^2)^{n+1}},$$

where $c_0 = -\int_0^1 \sigma^{2n-1} (1 + \sigma^2)^{2n} d\sigma$. The function $\phi(\tau, .)$ is not a Schwartz function on \mathbb{R}^3 , but we have that $(1 + \Delta)^k \phi(\tau, .) \in L^1(\mathbb{R}^3)$, hence $\widehat{(1 + |\xi|^2)^k \phi(\tau, .)}(\xi) \in L^\infty(\mathbb{R}^3)$. With these properties the inversion formula for the Radon transform (3.1) still holds. The proof follows straightforward from theorem 5.4 of [S-Sh]. Let us now compute the Radon transform of the function ϕ .

$$\begin{aligned} \mathcal{R}\phi(\tau, t, \xi) &= \int_{\mathbb{R}^2} \frac{16n}{\pi} \frac{4^{2n}(2n-1)!c_0}{(\tau^2 + 16(t^2 + u_1^2 + u_2^2))^{n+1}} du_1 du_2 = \\ &= \frac{16n}{\pi} \frac{4^{2n}(2n-1)!c_0}{16^{n+1}} \int_{\mathbb{R}^2} \frac{1}{\left(\frac{\tau^2}{16} + t^2 + (u_1^2 + u_2^2)\right)^{n+1}} du_1 du_2 = \\ &= \frac{16n}{\pi} \frac{4^{2n}(2n-1)!c_0}{16^{n+1}} \int_{-\frac{\pi}{2}}^{\frac{3}{2}\pi} \int_0^\infty \frac{\rho}{\left(\frac{\tau^2}{16} + t^2 + \rho^2\right)^{n+1}} d\rho d\theta = \\ &= \frac{4^{2n}(2n-1)!c_0}{(\tau^2 + 16|z|^2)^n}, \end{aligned}$$

where $z = t\xi$. Let $\varphi(\tau, z) = \frac{4^{2n}(2n-1)!c_0}{(\tau^2 + 16|z|^2)^n}$. Now from the expression of the fundamental solution of L in the classical case (see for example 4.3 of [G-S(2)]) we know that

$$\varphi(\tau, t\xi) = \sum_{k \geq 0} \frac{(-1)}{2k+2n} \int_{-\infty}^{\infty} e^{i\lambda t} L_k^{2n-1} \left(\frac{\lambda}{2} |\tau| \right) e^{-\frac{\lambda}{4} |\tau|} |\lambda|^{2n-1} d\lambda.$$

The first step to compute Φ is to change to polar coordinates in \mathbb{R}^3 the expression given in (1.4):

$$\begin{aligned} < \Phi, f > &= \sum_{k \in \mathbb{Z}_{\mathbb{R}^3}} \int \frac{1}{-|\lambda|(2k+2(p-q))} < \varphi_{w,k}, f > |w|^{2n} dw = \\ &= \sum_{k \in \mathbb{Z}_{S^2}} \int_0^\infty \int \frac{1}{-|\lambda|(2k+2(p-q))} < \varphi_{\lambda\xi,k}, f > |\lambda|^{2n+2} d\lambda d\xi. \end{aligned}$$

From the absolute convergence of (1.4) we can interchange the summation symbol with the integral

over S^2 . Because of $\Delta e^{i\lambda\langle\xi,z\rangle} = -|\lambda|^2 e^{i\lambda\langle\xi,z\rangle}$, we have that

$$\begin{aligned}
\langle \Phi, f \rangle &= \int_{S^2} \sum_{k \in \mathbb{Z}} \frac{(-1)}{(2k + 2(p - q))} \int_0^\infty \int_{N(p, q, \mathbb{H})} e^{i\lambda\langle\xi,z\rangle} \theta_{\lambda, k}(\alpha) f(\alpha, z) d\alpha dz |\lambda|^{2n+1} d\lambda d\xi = \\
&= \int_{S^2} \sum_{k \in \mathbb{Z}} \frac{1}{(2k + 2(p - q))} \int_0^\infty \int_{N(p, q, \mathbb{H})} \Delta e^{i\lambda\langle\xi,z\rangle} \theta_{\lambda, k}(\alpha) f(\alpha, z) d\alpha dz |\lambda|^{2n-1} d\lambda d\xi = \\
&= \int_{S^2} \sum_{k \in \mathbb{Z}} \frac{1}{(2k + 2(p - q))} \int_0^\infty \langle \varphi_{\lambda\xi, k}, \Delta f \rangle |\lambda|^{2n-1} d\lambda d\xi.
\end{aligned}$$

Next we break the summation indexes according to $k \geq 2q$, $k \leq -2p$ and $-2p < k < 2q$ to get the splitting $\langle \Phi, f \rangle = \langle \Phi_1, f \rangle + \langle \Phi_2, f \rangle$, and as in section we change summation index to get the series starting from $k = 0$. From the explicit definition of $\varphi_{\lambda\xi, k}$ we can write

$$\begin{aligned}
\langle \Phi_1, f \rangle &= \int_{S^2} \sum_{k \geq 0} \frac{1}{2k + 2n} \int_0^\infty \int_{\mathbb{R}^3} e^{i\lambda\langle\xi,z\rangle} \times \\
&\quad \times \langle T_{\lambda, k+2q} - T_{\lambda, -k-2p}, N\Delta f(., z) \rangle dz |\lambda|^{2n-1} d\lambda d\xi,
\end{aligned}$$

where $T_{\lambda, k}$ is defined by equations (1.6) and (1.7). By performing similar computations than in section 2 and introducing the function

$$G_f(\tau, z) = Nf(\tau, z) - \sum_{j=0}^{2n-2} \frac{\partial^j Nf}{\partial \tau^j}(0, z) \frac{\tau^j}{j!}$$

we get the splitting $\langle \Phi_1, f \rangle = \langle \Phi_{11}, f \rangle + \langle \Phi_{12}, f \rangle$, where

$$\begin{aligned}
\langle \Phi_{11}, f \rangle &= \int_{S^2} \sum_{k \geq 0} \frac{(-1)}{2k + 2n} \int_0^\infty \int_{\mathbb{R}^3} \int_{-\infty}^\infty e^{i\lambda\langle\xi,z\rangle} \times \\
&\quad \times \operatorname{sgn}(\tau) L_k^{2n-1} \left(\frac{2}{\lambda} |\tau| \right) e^{-\frac{\lambda}{4} |\tau|} \Delta G_f(\tau, z) d\tau dz |\lambda|^{2n-1} d\lambda d\xi, \tag{3.2}
\end{aligned}$$

$$\begin{aligned}
\langle \Phi_{12}, f \rangle &= 2 \int_{S^2} \sum_{k \geq 0} \frac{1}{2k + 2n} \int_0^\infty \int_{\mathbb{R}^3} e^{i\lambda\langle\xi,z\rangle} \times \\
&\quad \times \sum_{\substack{l=0 \\ l \text{ odd}}}^{2n-2} \left(\frac{2}{\lambda} \right)^{l+1} (a_{kl} + b_{kl}) \langle \delta^{(l)}, \Delta Nf(., z) \rangle dz |\lambda|^{2n-1} d\lambda d\xi, \tag{3.3}
\end{aligned}$$

and a_{kl} , b_{kl} are the same constant defined in (2.8) and (2.3), respectively. Now let us recall the fact that

$$\int_{S^2} \int_0^\infty e^{i\lambda\langle\xi,z\rangle} F(|\lambda|) d\lambda d\xi = \frac{1}{2} \int_{S^2} \int_{-\infty}^\infty e^{i\lambda\langle\xi,z\rangle} F(|\lambda|) d\lambda d\xi,$$

and we apply the dual Radon transform to (3.2).

Observe now that

$$\int_{-\infty}^{\infty} \int_{\mathbb{R}^3} \frac{\operatorname{sgn}(\tau) G_f(\tau, z)}{(1 + 16|z|^2)^{n+1}} dz d\tau$$

converges, which we can see by changing to polar variables in \mathbb{R}^3 and arguing like in lemma 2.2 of [G-S(2)].

We finally get that

$$\begin{aligned} \langle \Phi_{11}, f \rangle &= \frac{1}{2} \langle -2\pi \frac{16n}{\pi} \frac{4^{2n}(2n-1)!c_0}{(\tau^2 + 16|z|^2)^{n+1}}, \operatorname{sgn}(\tau) G_f(\tau, z) \rangle = \\ &= -4^{2n+2} n(2n-1)!c_0 \langle \frac{1}{(\tau^2 + 16|z|^2)^{n+1}}, \operatorname{sgn}(\tau) G_f(\tau, z) \rangle. \end{aligned}$$

We have thus proven that the expression defining Φ_{11} is finite. Then the expression defining Φ_{12} is also finite, and by Abel's lemma we can write

$$\langle \Phi_{12}, f \rangle = 2 \lim_{r \rightarrow 1^-} \lim_{\epsilon \rightarrow 0^+} \sum_{k \geq 0} \sum_{\substack{l=0 \\ l \text{ odd}}}^{2n-2} 2^{l+1} (a_{kl} + b_{kl}) \frac{r^{2k+2n}}{2k+2n} d_{\epsilon,l,f},$$

where

$$d_{\epsilon,l,f} = \int_{S^2} \int_0^\infty \int_{\mathbb{R}^3} e^{-\epsilon\lambda} e^{i\lambda \langle \xi, z \rangle} |\lambda|^{2n-l-2} \langle \delta^{(l)}, \Delta N f(., z) \rangle dz d\lambda d\xi. \quad (3.4)$$

We need to compute $\lim_{\epsilon \rightarrow 0^+} d_{\epsilon,l,f}$. Observing that $\Delta e^{i\lambda \langle \xi, z \rangle} = -|\lambda|^2 e^{i\lambda \langle \xi, z \rangle}$, we have that

$$\begin{aligned} d_{\epsilon,l,f} &= (-1)^{l+1} \int_{S^2} \int_0^\infty \int_{\mathbb{R}^3} e^{-\epsilon\lambda} e^{i\lambda \langle \xi, z \rangle} |\lambda|^{2n-l} \frac{\partial^l}{\partial \tau^l} N f(0, z) dz d\lambda d\xi = \\ &= (-1)^{l+1} \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} e^{-\epsilon|x|} |x|^{2n-l-2} e^{i\langle x, z \rangle} \frac{\partial^l}{\partial \tau^l} N f(0, z) dz dx, \end{aligned}$$

where we have changed to cartesian coordinates in \mathbb{R}^3 . To solve this integral let us observe that

$$(-1)^{2n-l-2} e^{-\epsilon|x|} |x|^{2n-l-2} = \widehat{\frac{\partial^{2n-l-2}}{\partial \epsilon^{2n-l-2}}} P_\epsilon(x),$$

where $P_\epsilon(x)$ is the Poisson kernel and $\widehat{\cdot}$ denotes the Fourier transform. Let us write

$$d_{\epsilon,l,f} = (-1)^l \int_{\mathbb{R}^3} \widehat{\frac{\partial^{2n-l-2}}{\partial \epsilon^{2n-l-2}}} P_\epsilon(x) \left(\widehat{\frac{\partial^l}{\partial \tau^l} N f(0, .)} \right) (x) dx = (-1)^l \frac{\partial^{2n-l-2}}{\partial \epsilon^{2n-l-2}} (P_\epsilon * h)(0).$$

Taking limit as $\epsilon \rightarrow 0^+$ we obtain

$$\lim_{\epsilon \rightarrow 0^+} = (-1) (-\Delta)^{\frac{2n-l-2}{2}} \frac{\partial^l}{\partial \tau^l} N f(0, 0),$$

where $(-\Delta)^{\frac{2n-l-2}{2}}$ is a fractionary exponential of the Laplacian (see for example [S-Sh]), which is the operator defined for $g \in \mathcal{S}(\mathbb{R}^3)$ by

$$(-\Delta)^{\frac{2n-l-2}{2}}g(x) = \int_{\mathbb{R}^3} |\omega|^{2n-l-2} \hat{g}(\omega) e^{i\langle \omega, z \rangle} d\omega.$$

This computation together with proposition 4.8 of [G-S(2)] lets us write

$$\langle \Phi_{12}, f \rangle = \sum_{\substack{l=0 \\ l \text{ odd}}}^{2n-2} \sum_{j=1}^{l+1} \frac{1}{2^{2n-2l-j-3}} c_j \binom{2n-j-1}{l-j-1} (-1) (-\Delta)^{\frac{2n-l-2}{2}} \frac{\partial^l}{\partial \tau^l} Nf(0,0),$$

where each c_j is the constant defined in Remark 4.7 of [G-S(2)].

After performing the usual computations for Φ_2 we have proved the main theorem of this section:

Theorem 3.1. *Let c_0 and c_j be the constants defined above. Then*

$$\begin{aligned} \langle \Phi, f \rangle &= -4^{2n+2} n(2n-1)! c_0 \langle \frac{1}{(\tau^2 + 16|z|^2)^{n+1}}, \operatorname{sgn}(\tau) G_f(\tau, z) \rangle + \\ &+ \sum_{\substack{l=0 \\ l \text{ odd}}}^{2n-2} \sum_{j=1}^{l+1} \frac{1}{2^{2n-2l-j-3}} c_j \binom{2n-j-1}{l-j-1} (-1) (-\Delta)^{\frac{2n-l-2}{2}} \frac{\partial^l}{\partial \tau^l} Nf(0,0) + \\ &+ \sum_{-2q < k < 0} \frac{1}{2k+2(p-q)} \left[(-1)^{k+1} \sum_{r=0}^{k+2p-1} \binom{k+2p-1}{r} 2^{-k-2p+2r+2} \times \right. \\ &\quad \times (-1) (-\Delta)^{\frac{2n-r-2}{2}} \frac{\partial^r}{\partial \tau^r} Nf(0,0) + \\ &+ \sum_{l=0}^{2n-2-k-2p} (-1)^{-l-k} \binom{-k-2p+2n-1}{2n-2-l-k-2p} \sum_{r=0}^{l+k-2p} \binom{l+k-2p}{r} \times \\ &\quad \times (-1) 2^{-l-k+2p+2r+1} (-\Delta)^{\frac{2n-r-2}{2}} \frac{\partial^r}{\partial \tau^r} Nf(0,0) \Big] + \\ &+ \sum_{0 \leq k < 2q} \frac{1}{2k+2(p-q)} \left[(-1)^{k+1} \sum_{r=0}^{-k+2q-1} \binom{-k+2q-1}{r} \times \right. \\ &\quad \times 2^{k-2q+2r+1} (-1)^r (-\Delta)^{\frac{2n-r-2}{2}} \frac{\partial^r}{\partial \tau^r} Nf(0,0) + \\ &+ \sum_{l=0}^{2n-2+k-2q} (-1)^{-l+k-2q} \binom{k-2q+2n-1}{2n-2-l+k-2q} \sum_{r=0}^{l-k+2q} \binom{l-k+2q}{r} \times \\ &\quad \times 2^{-l+k-2q+2r+1} (-1)^r (-\Delta)^{\frac{2n-r-2}{2}} \frac{\partial^r}{\partial \tau^r} Nf(0,0) \Big]. \end{aligned}$$

Remark

Let N be a group of Heisenberg type and let η be its Lie algebra. So $\eta = V \oplus \mathfrak{z}$, with $\dim V = 2m$ and $\dim \mathfrak{z} = k$. Let $U(V)$ be the unitary group acting on V . Then it is known ([R]) that $(N \ltimes U(V), U(V))$

is a Gelfand pair, and also in [R] were computed the spherical functions. We fix an orthonormal basis of V , $\{X_1, \dots, X_{2m}\}$, and consider the operator

$$L = \sum_{j=1}^{2m} X_j^2.$$

With the same arguments as above, using the Radon transform in \mathbb{R}^k and the fundamental solution of L in the classical Heisenberg group $2m + 1$ dimensional, we can recover the fundamental solution of L (see [K], [R]).

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